

# Efficient Electromagnetic Modeling of Interconnects and Packages in the dc-to-Skin Effect Limit

---

**Dean Neikirk\***

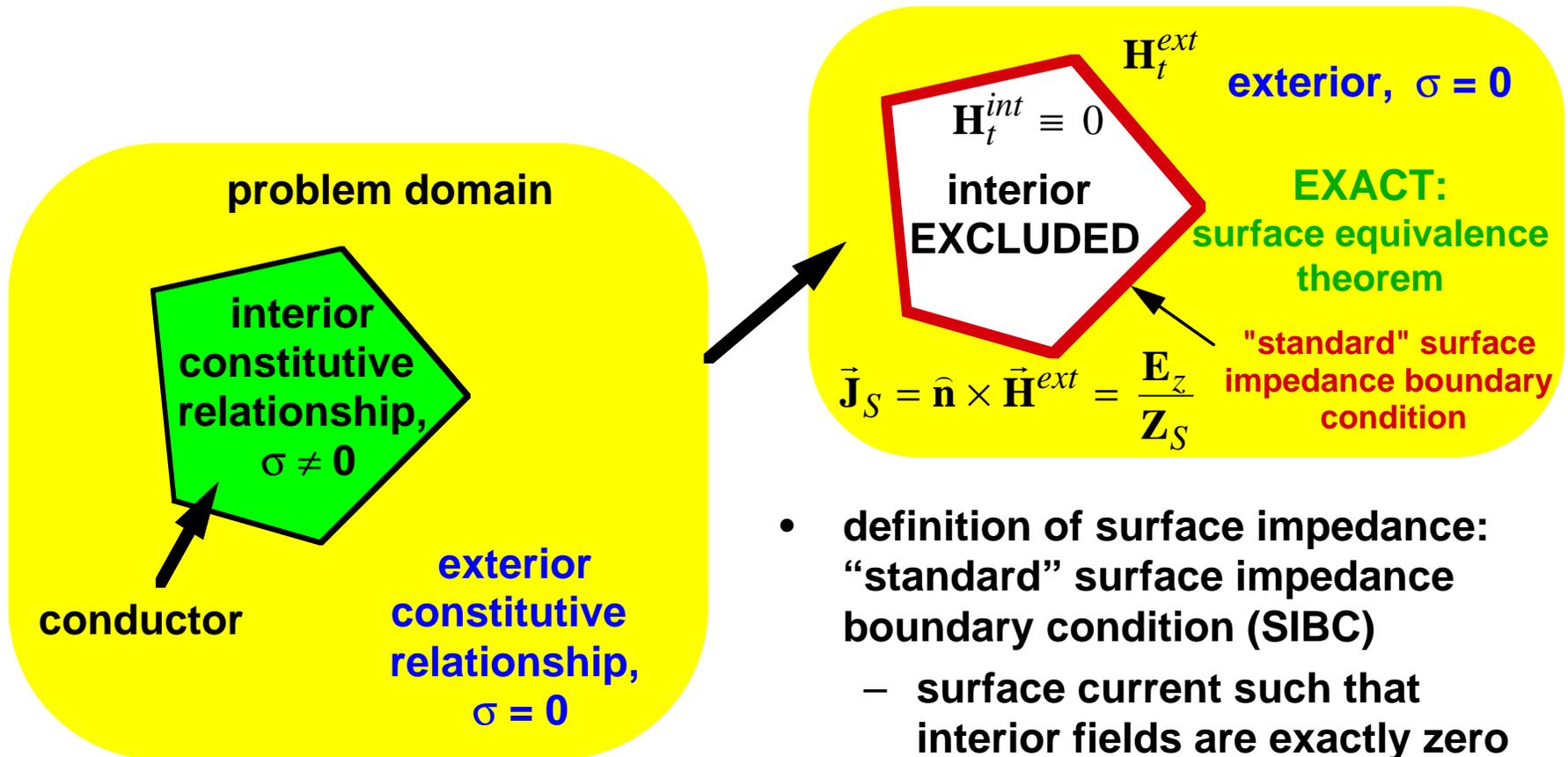
**Guanghan Xu, and Leszek Demkowicz**

**Department of Electrical and Computer Engineering  
The University of Texas at Austin**

**DARPA PM: Dr. Anna Tsao (formerly Dr. Nicholas Naclerio)  
Electronic Packaging and Interconnects Program  
Contracting Agent: Dr. Arje Nachman, AFOSR**

\* e-mail: [neikirk@mail.utexas.edu](mailto:neikirk@mail.utexas.edu)  
www: <http://weewave.mer.utexas.edu>  
phone: 512-471-4669

# “Efficient” modeling: Elimination of conductors from solution domain via surface impedance



- definition of surface impedance: “standard” surface impedance boundary condition (SIBC)
  - surface current such that interior fields are exactly zero

# Utility of surface impedance BCs

---

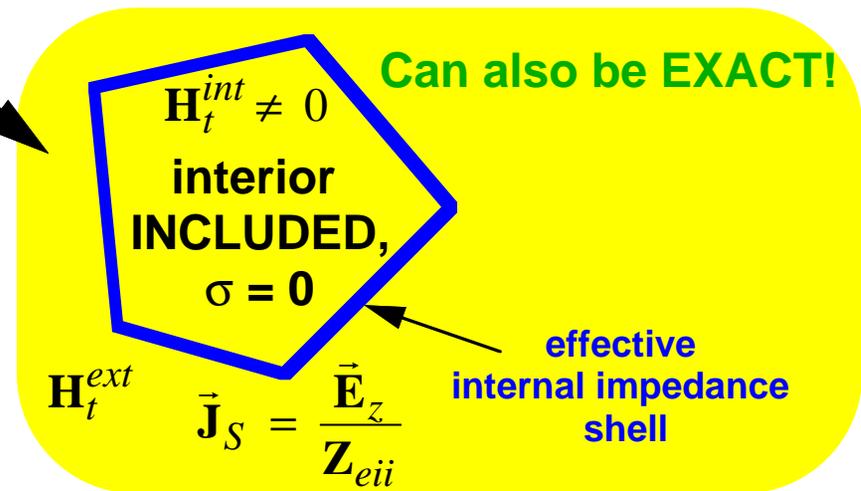
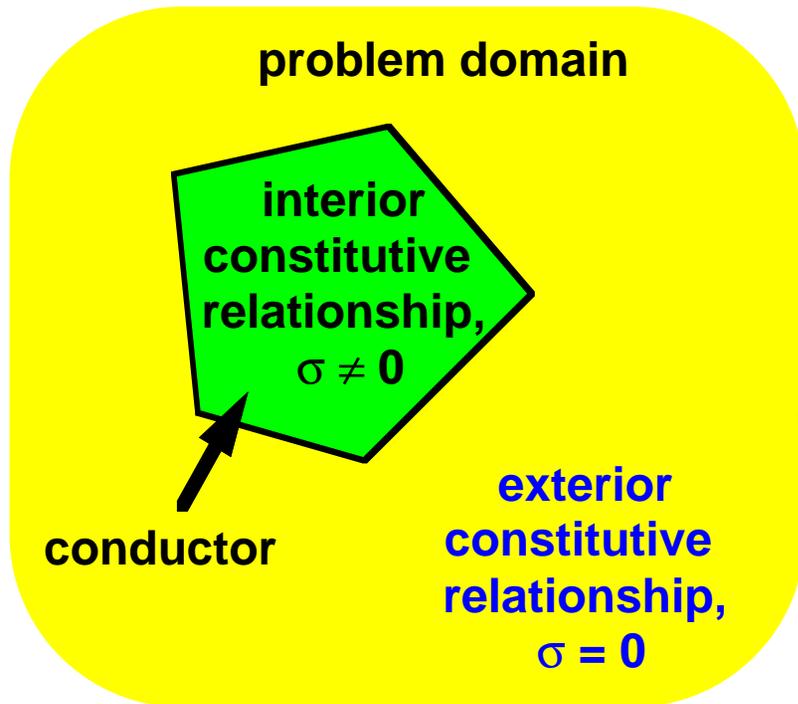
- simplifies problem only if known *a priori*
  - how well can the boundary condition be approximated?
  - most common approximation for SIBC:

$$Z_s = \frac{(1 + j)}{\sigma \cdot \delta}$$

- valid only at high frequency where E/H is a local property
- requires conductor dimensions  $\gg$  skin depth
- fails at “low” frequency and when line-to-line coupling is strong

# Alternative definition of conductor surface impedance: EII

- “effective” internal impedance boundary condition (EII)
  - interior of conductor replaced with non-conducting exterior medium
  - interior magnetic field is **not** identically zero



# Alternative definition of conductor surface impedance: EII

- “effective” internal impedance boundary condition (EII)

$$Z_{eii}(\omega, r') = \frac{E_z(r')}{H_t^{ext}(r') - H_t^{in}(r')} = \frac{E_z(r')}{J_s(r')}$$

- $E_z$  and  $H_{t,ext}$  from the solution of

$$\sum_{q=1}^m \oint_{\Gamma_q} dr' G_{ext}^1(r, r') j\omega\mu H_t^{ext}(r') - \sum_{q=1}^m \oint_{\Gamma_q} dr' \left[ G_{ext}^2(r, r') - 0.5\delta(r - r') \right] \cdot \left( E_z(r') + \nabla_z \Phi^q \right) = 0$$

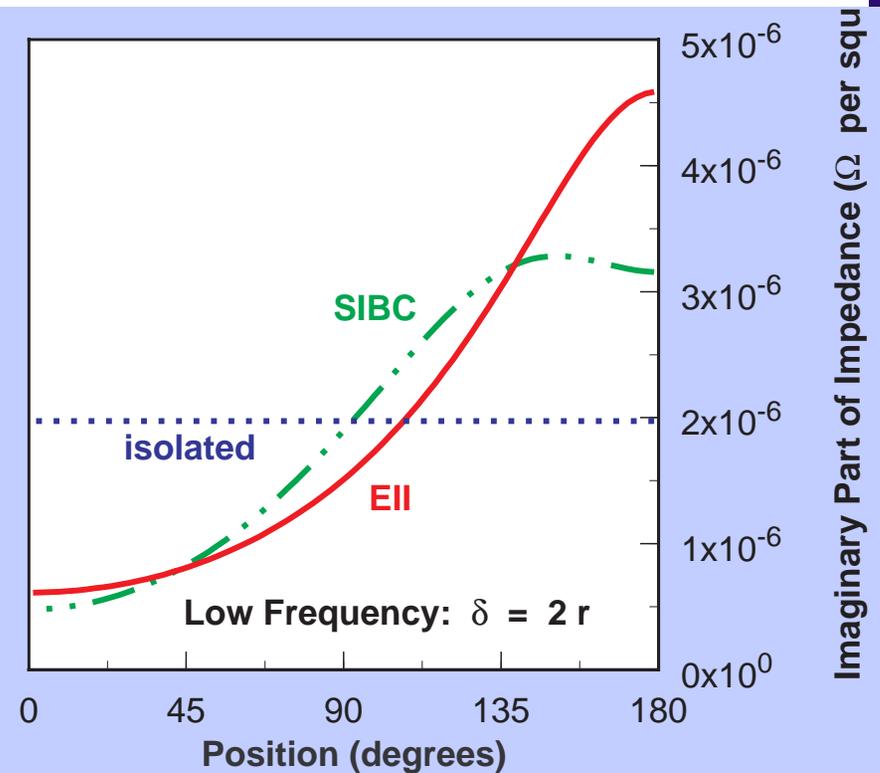
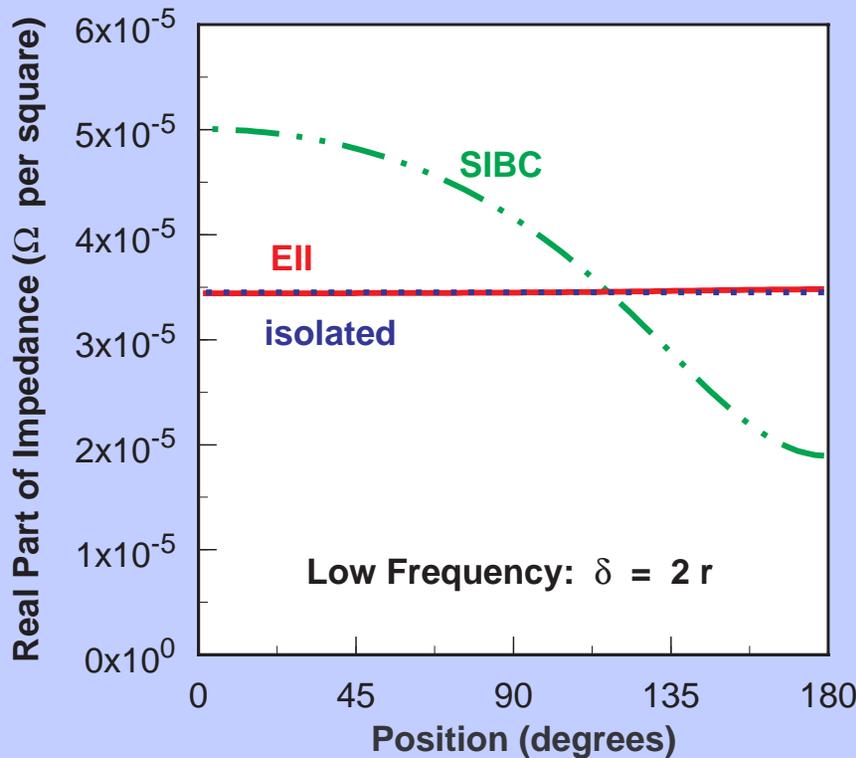
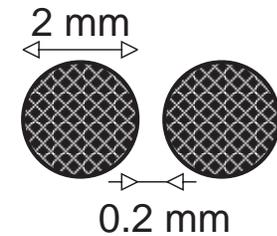
- $H_{t,in}$  from the solution of

$$\oint_{\Gamma_q} dr' G_q^1(r, r') j\omega\mu H_t^{in}(r') - \oint_{\Gamma_q} dr' \left[ G_q^2(r, r') + 0.5\delta(r - r') \right] E_z(r') = 0$$

- utility of impedance boundary condition determined by how easily it can be approximated
  - isolated conductor impedance is easier to estimate than coupled line surface impedance

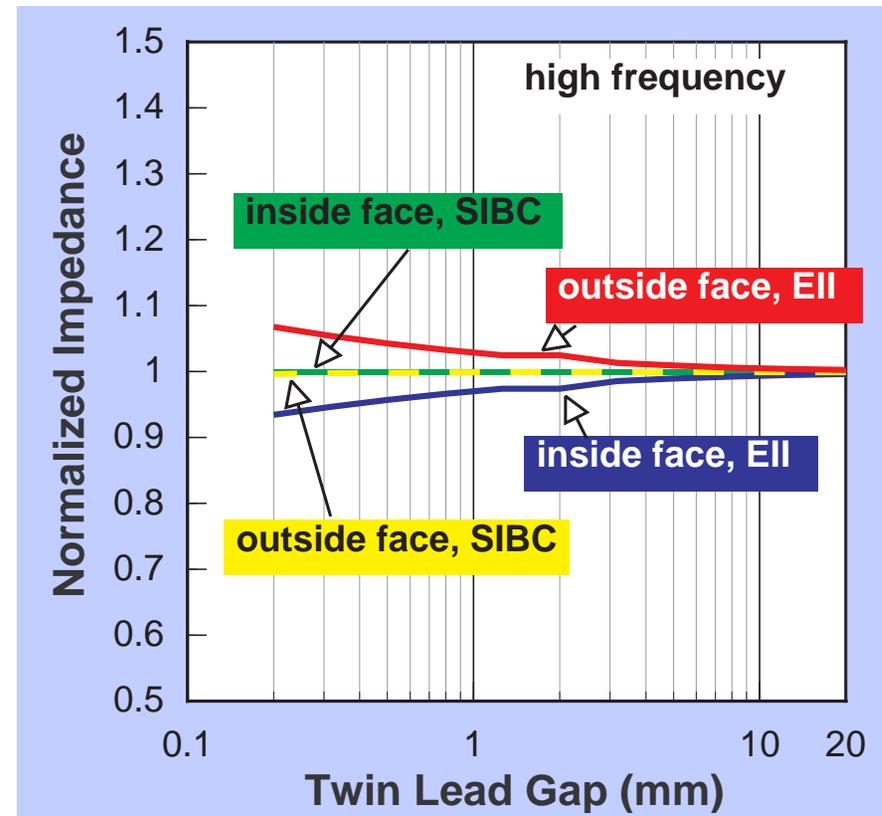
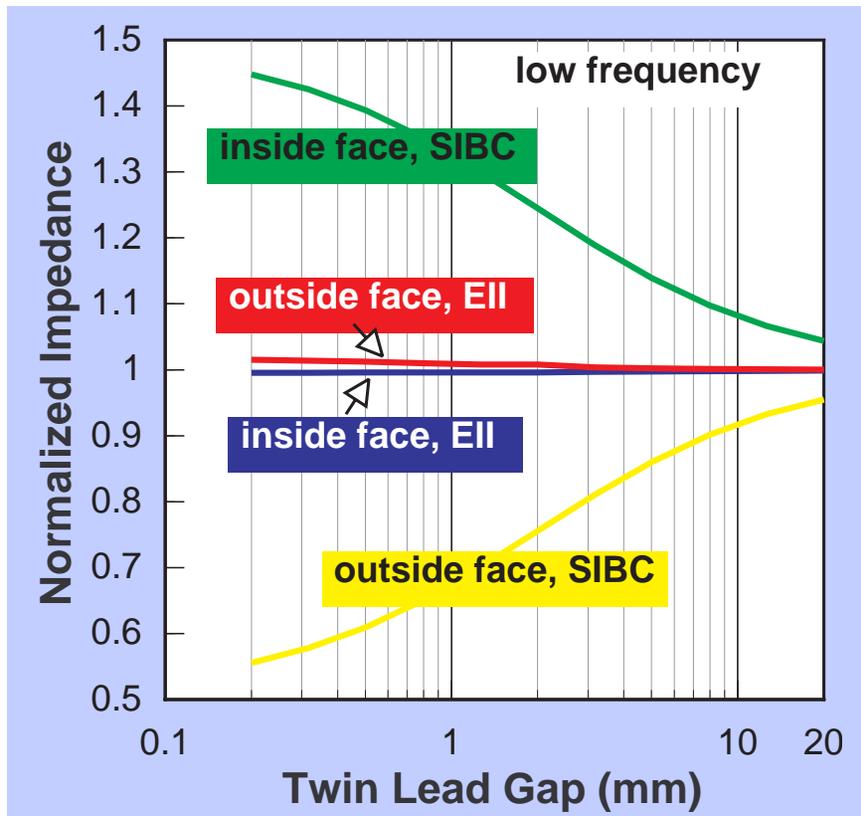
# Surface impedance comparisons

- simple circular twin lead
- all tend to  $1/\sigma\delta$  at high frequency
- EII well approximated by isolated conductor surface impedance at low frequency

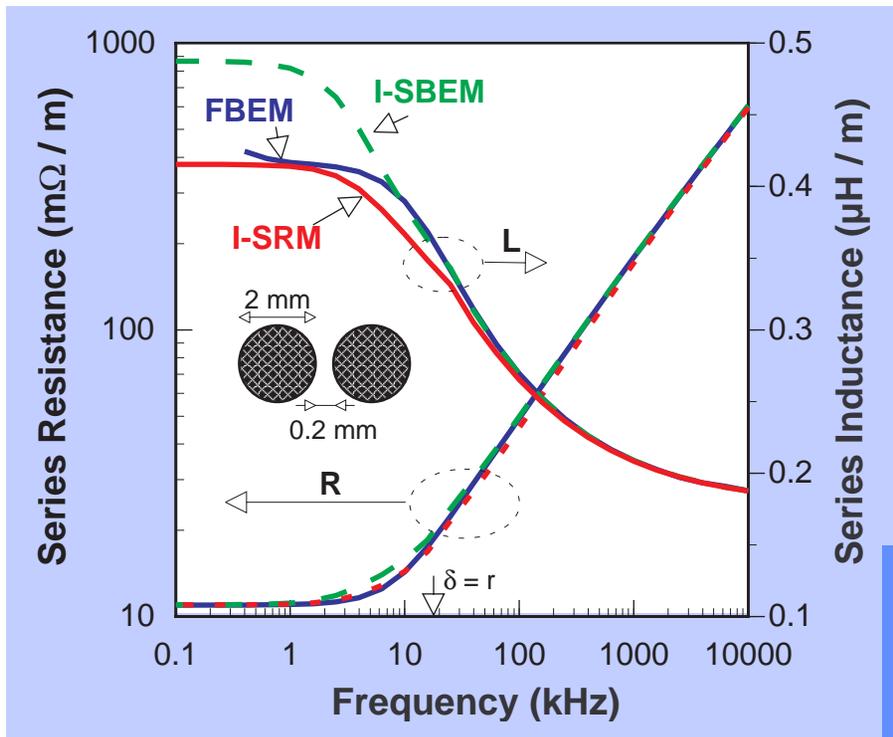


# Dependence on coupling

- normalized to magnitude of isolated conductor surface impedance
- 2 mm radius circular conductors



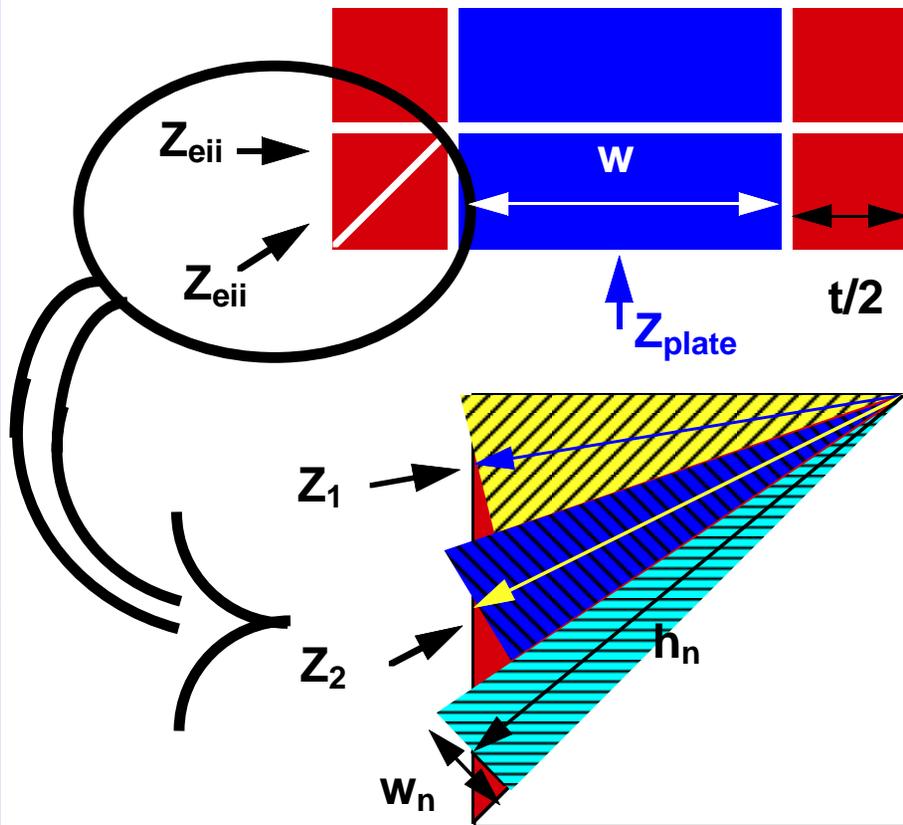
# Circular twin lead comparisons



- techniques used
  - boundary element method (BEM)
  - “surface ribbon” technique (SRM)

Method	number of unknowns	CPU time (assemble)	CPU time (solve)
BEM	573	468	100
BEM with SIBC	290	18	8
SRM with EII	287	5	9

# EII approximations for rectangular conductors : "Transmission line" model



- decompose bar into rectangular and square sections
- central rectangular regions: simple skin depth in "flat" plate

wide plate approximation:

$$Z_{plate} = \frac{(1 + j) / (\sigma \cdot \delta)}{\tanh[(1 + j) \cdot t / (2 \cdot \delta)]} \cdot \frac{1}{w}$$

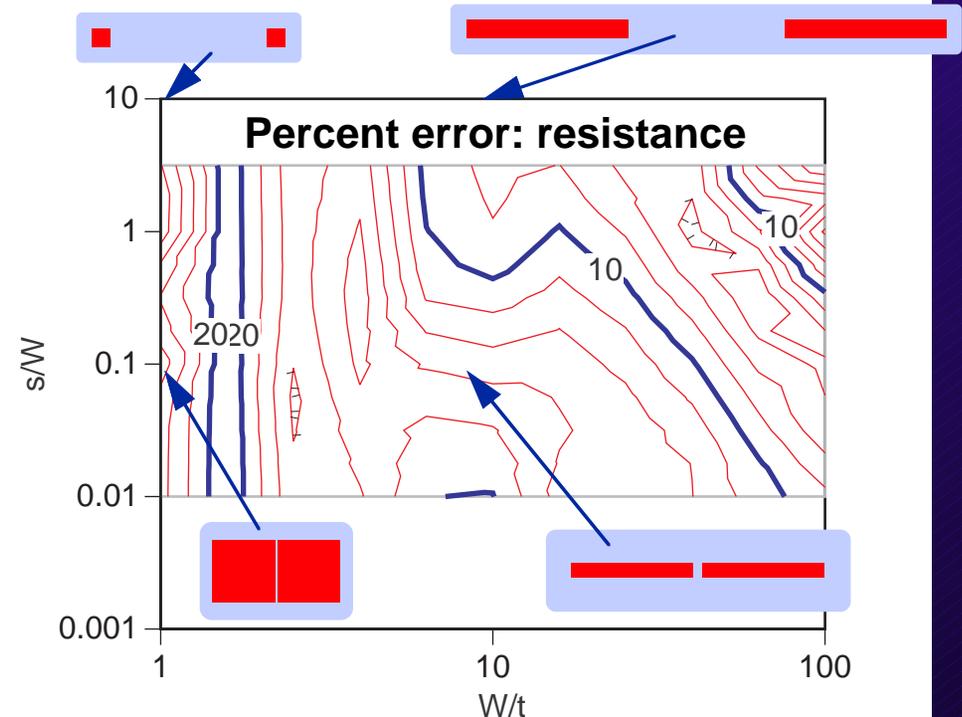
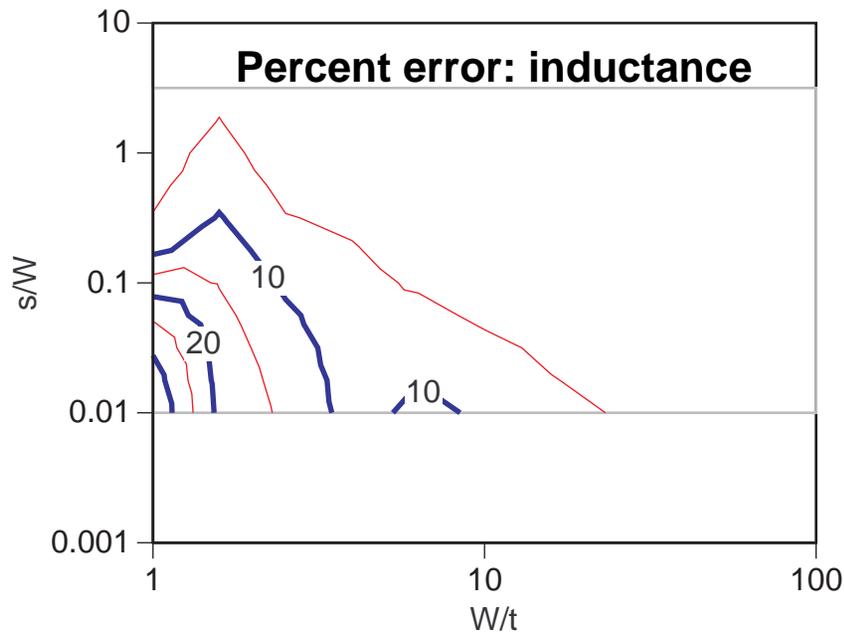
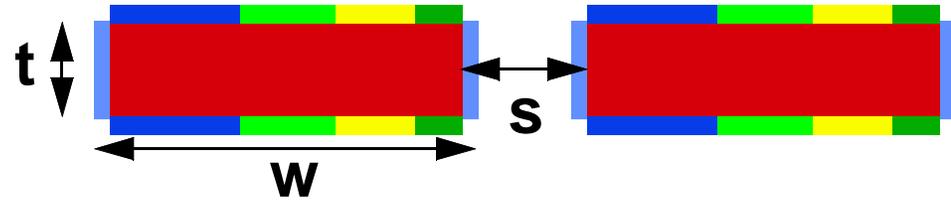
- corner regions: decompose t/2 corner into N triangles
  - exact dc resistance in limit of N large
  - high frequency behavior
    - effective thickness > t/2, current crowds toward "corner" at lower frequency

$$Z_n = \frac{j \cdot (1 + j)}{\sigma \cdot \delta} \cdot \frac{J_0(j \cdot [1 + j] \cdot h_n / \delta)}{J_1(j \cdot [1 + j] \cdot h_n / \delta)} \cdot \frac{1}{w_n}$$

# Surface ribbon method for rectangular "twin lead"

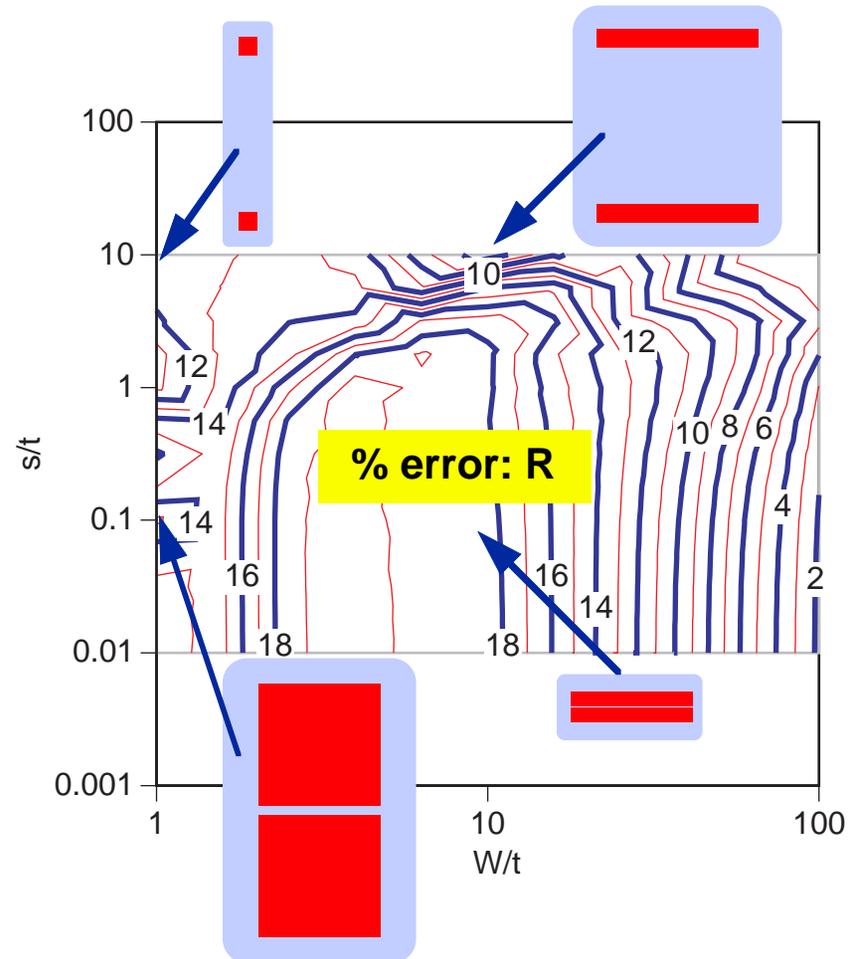
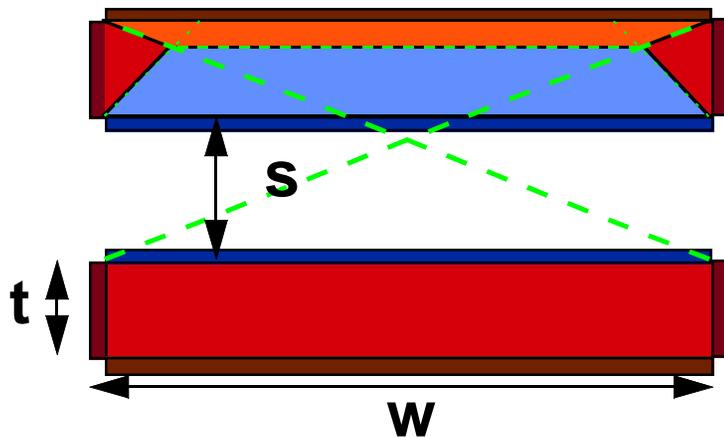
## minimum ribbons:

- at most five ribbons used for "wide" faces
- one ribbon used for "narrow" faces



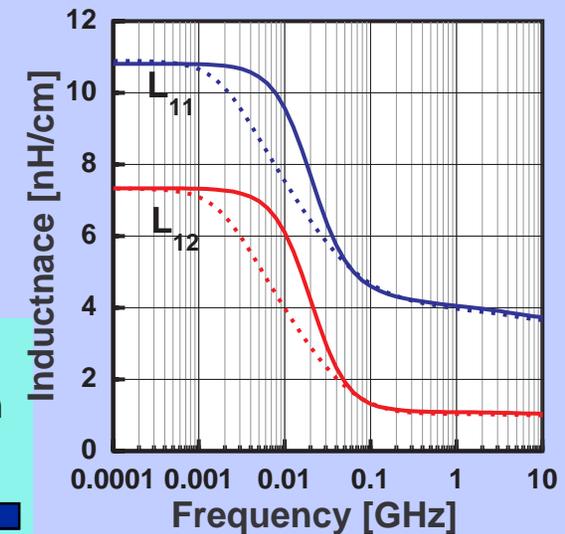
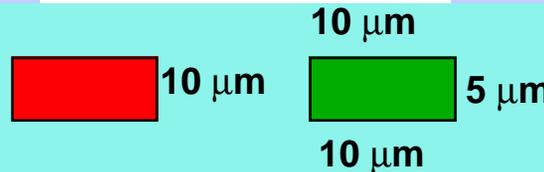
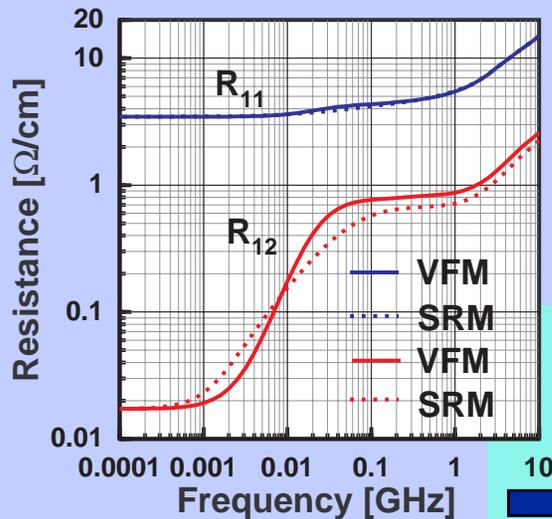
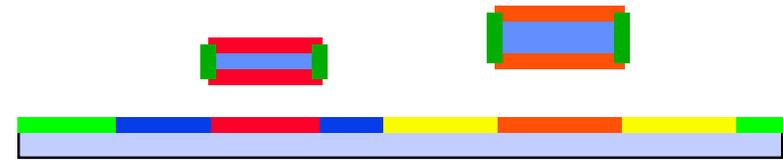
# Parallel plate errors using minimum surface ribbons

- only one ribbon per face required with appropriate EII model

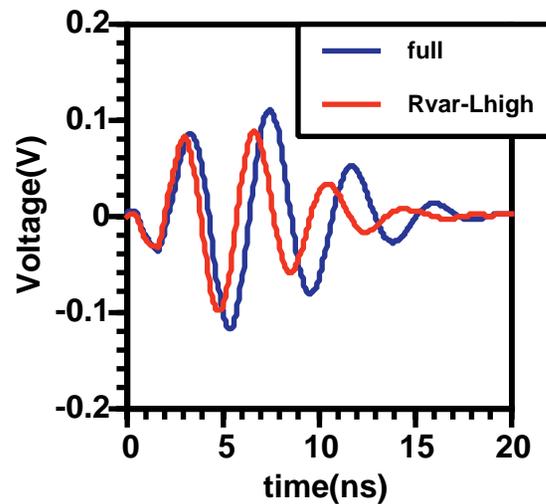
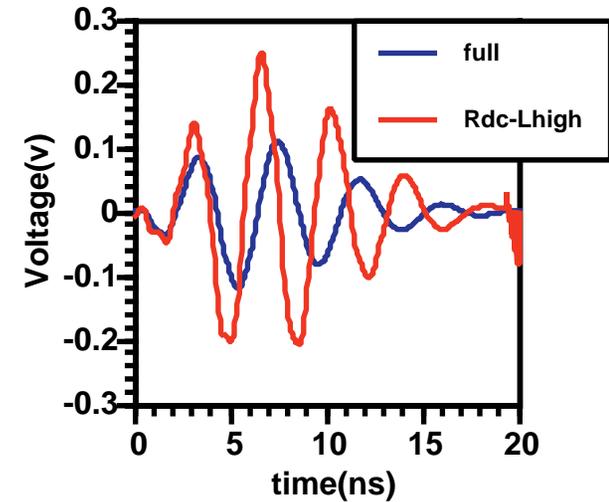
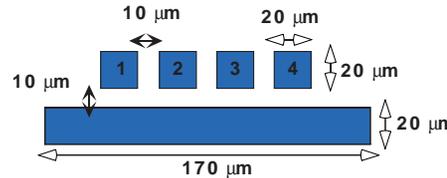
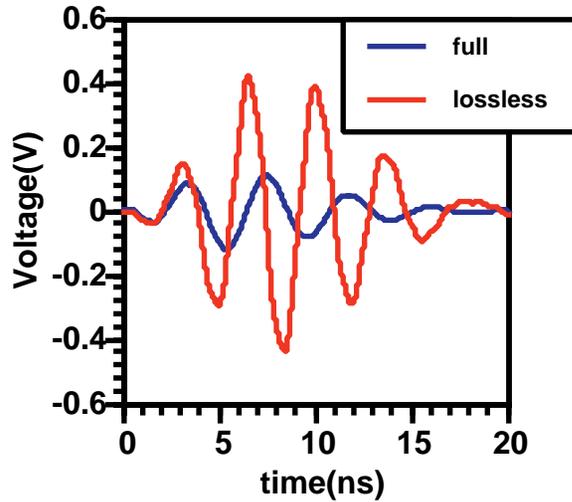


# Applications: Minimum ribbon segmentation for coupled microstrip over finite ground plane

- signal lines : **4** segments each signal line (one per face)
- ground plane : **7 ~ 9** segments
- unknowns - VFM : **492** ; SRM : **15**
- run time - VFM : **39 sec**; SRM : **0.023 sec**



# Effects of the frequency dependencies on time domain waveforms



- line 1 excited with 0.1ns rise and fall time
- measured far end of line 4 (length=0.1 m)
- $R_S = 5 \text{ Ohm}$ ,  $C_L = 10 \text{ pF}$

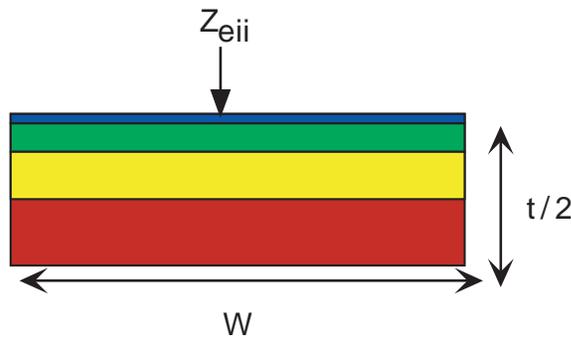
• must include both R and L dependence on frequency / time

## Time domain simulation using surface impedance BCs

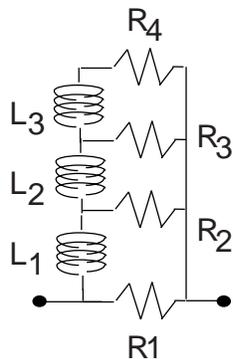
---

- **conventional approach: convert surface impedance from frequency domain to time domain**
  - can use Prony's method or Chebyshev approximation to find time domain exponential form
    - tends to produce many terms
- **alternative: use compact ladder circuit to formulate time domain surface impedance boundary condition**
  - use circuit to determine time domain exponential form
  - leads to time domain formulation very similar to lossless result
    - matrix containing finite conductivity effects is time independent, requires inversion only **once**

# Equivalent circuit modeling for EII: transforming frequency domain EII into time domain



- each “isolated conductor” cross-section divided into 4 parts
- each part represented with 1 resistor and 1 inductor
- rules to determine values of circuit elements:



$$\frac{R_1}{R_2} = \frac{R_2}{R_3} = \frac{R_3}{R_4} = RR \quad \frac{L_1}{L_2} = \frac{L_2}{L_3} = LL$$

- additional constraints: correct DC resistance and inductance

- RR and LL are empirically determined constants unique to the geometry of the conductor

$$Z_{eii}(s) = R_1 \frac{b_3 s^3 + b_2 s^2 + b_1 s + b_0}{a_3 s^3 + a_2 s^2 + a_1 s + a_0}$$

# Compact circuit model as a replacement for the EII

- constraints generate:

$$RR^3 + RR^2 + RR + (1 - \alpha_R) = 0$$

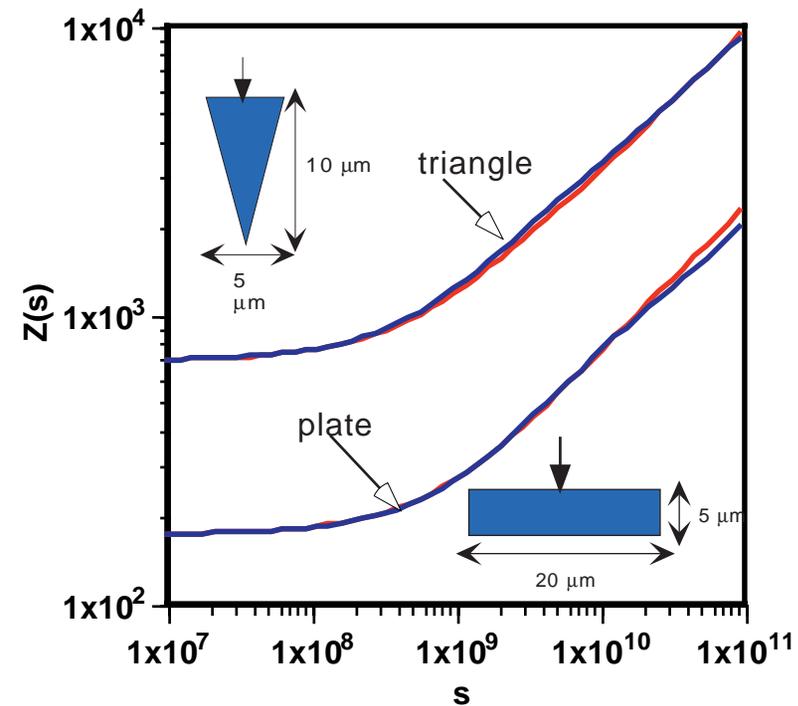
$$\left(\frac{1}{LL}\right)^2 + \left(1 + \frac{1}{RR}\right)^2 \frac{1}{LL} + \left(\left[\frac{1}{RR}\right]^2 + \frac{1}{RR} + 1\right)^2 - \alpha_L \left(\left[1 + \frac{1}{RR}\right] \left[\left\{\frac{1}{RR}\right\}^2 + 1\right]\right)^2 = 0$$

$$\alpha_R = C_1 \frac{p}{\delta_{\max}} = \frac{R_1}{R_{dc}} \quad \delta_{\max} = \sqrt{\frac{2}{\omega_{\max} \mu_o \sigma}} \quad \alpha_L = C_2 \alpha_R = L_{dc} / L_1$$

- p** is the conductor depth parameter, **C<sub>1</sub>** and **C<sub>2</sub>** depend on geometry:
  - triangle for corners: **p** = height, **C<sub>1</sub>** = 0.56, **C<sub>2</sub>** = 0.315
  - half plate for mid regions: **p** = thickness, **C<sub>1</sub>** = 10.8, **C<sub>2</sub>** = 0.2
  - circular conductors: **p** = **r**, **C<sub>1</sub>** = 0.53, **C<sub>2</sub>** = 0.315
- ladder values completely determined for each ribbon on conductor surface

# Time domain conversion using equivalent circuit model

- equivalent circuit model
  - can be easily constructed
  - rational function in s-domain, exponential function in time domain
  - problem size can be reduced using Pade approximation: dominant pole reduction
  - time domain convolution problem can be solved using recursive properties



blue: numerical result

red: circuit model

# Derivation of time domain equation

- frequency domain (s-domain) equation

$$\left[ \frac{Z_{eii}}{s} \right] s[I] + [L]s[I] = -\frac{\partial}{\partial z}[V]$$

- transformation into time domain

$$L^{-1}\left(\left[ \frac{Z_{eii}}{s} \right] = \left[ R_1 \frac{b_3 s^3 + b_2 s^2 + b_1 s + b_0}{s(a_3 s^3 + a_2 s^2 + a_1 s + a_0)} \right] \right)$$

$$= [R_1 \sum K_i e^{p_i t}] = [\zeta(t)]$$

$$[\zeta(t)] * \frac{\partial}{\partial t}[I] + [L] \frac{\partial}{\partial t}[I] = -\frac{\partial}{\partial z}[V]$$

- time domain convolution

$$Y(n\Delta t) = X(n\Delta t) * Ke^{p(n\Delta t)}$$

$$= K\Delta t \cdot X(n\Delta t) + e^{p\Delta t} \cdot Y((n-1)\Delta t)$$

- application of recursive equation

$$[K] \frac{\partial}{\partial t}[I] + [L] \frac{\partial}{\partial t}[I] + [V_{ds}] = -\frac{\partial}{\partial z}[V]$$

$$[L'] \frac{\partial}{\partial t}[I] + [V_{ds}] = -\frac{\partial}{\partial z}[V]$$

$$[C] \frac{\partial}{\partial t}[V] + [G] \cdot [V] = -\frac{\partial}{\partial z}[I]$$

- lossless like equation with additional voltage source
- voltage source depends on poles, residues, time step, and values from previous time step
- different simulators can be used to solve equations

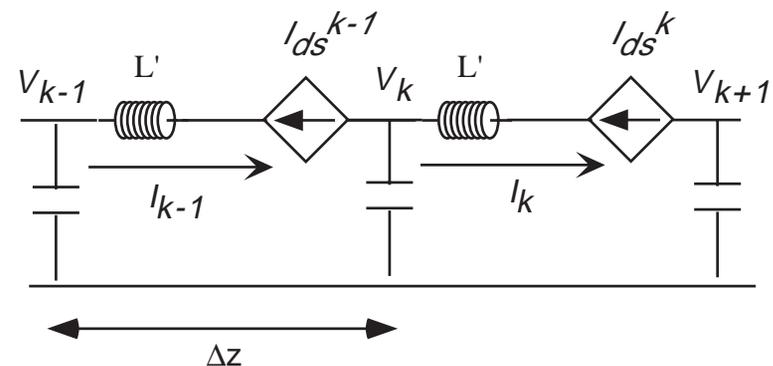
## Example I: time stepping solution (FDTD)

- make finite difference approximation to the partial derivatives: extra current source compared to lossless case
- each voltage and adjacent current solution point separated by  $\Delta z/2$
- $\Delta t$  has to be kept small to satisfy stability condition: may not be appropriate for electrically long lines

$$[I]_k^{n+\frac{2}{3}} = [I]_k^{n+\frac{1}{2}} - \left[ \frac{L'}{\Delta t} \right]^{-1} \left( \frac{[V]_{k+1}^{n+1} - [V]_k^{n+1}}{\Delta z} + [V_{ds}]_k^{n+1} \right)$$

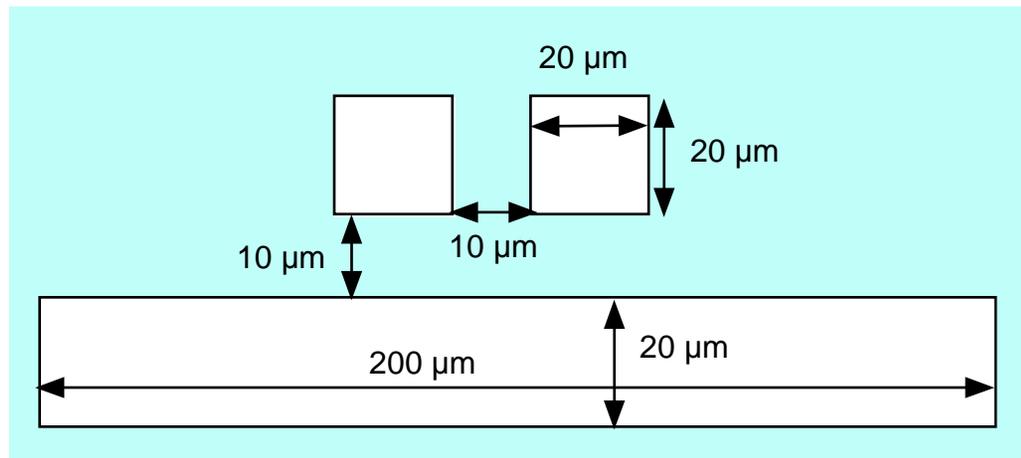
$$= [I]_k^{n+\frac{1}{2}} - \left[ \frac{L'}{\Delta t} \right]^{-1} \left( \frac{[V]_{k+1}^{n+1} - [V]_k^{n+1}}{\Delta z} \right) - [I_{ds}]_k^{n+1}$$

$$[V]_k^{n+1} = [V]_k^n - \left[ \frac{C}{\Delta t} \right]^{-1} \left( \frac{[I]_{k+1}^{n+\frac{1}{2}} - [I]_k^{n+\frac{1}{2}}}{\Delta z} \right)$$



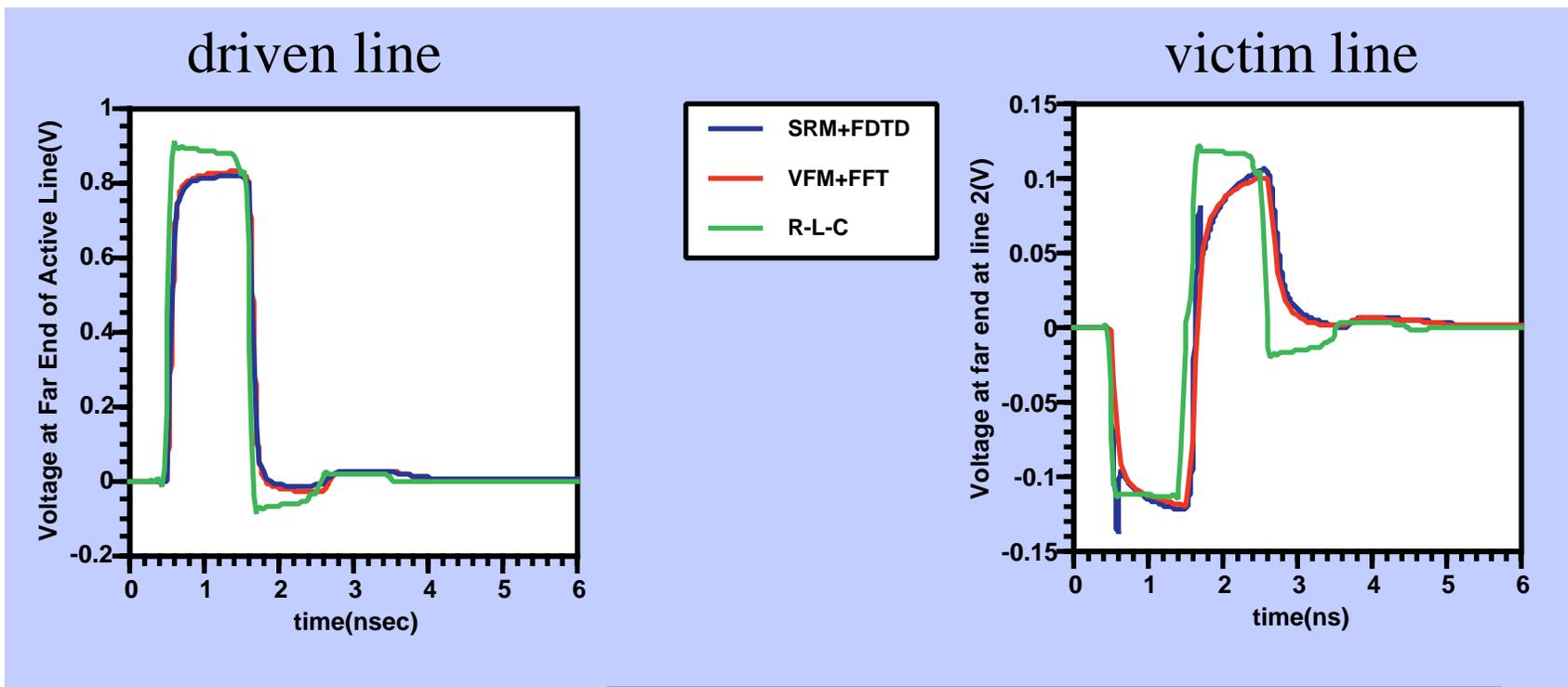
# EI/Circuit approximation in FDTD

- test case geometry: two  $20\mu\text{m}$  square microstrip lines,  $30\mu\text{m}$  pitch, all conductivities finite



- $10\text{cm}$  length,  $5\Omega$  source impedance,  $50\Omega$  load
- comparisons
  - finite difference time domain method using EI circuit derived BC
  - FFT using frequency domain dispersion curve
    - conventional current filament method (FM)
    - surface ribbon method (SRM)
  - simple RLC transmission line model (no skin effect)

# Driven and victim line comparisons

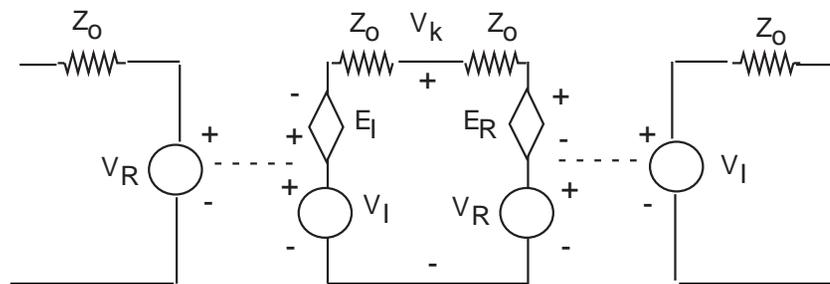


- single input pulse, 0.1nsec rise & fall times, 1nsec on time

method	pre-calculation	main calculation (FDTD or FFT)
SRM-FDTD	0.5 sec	89.5 sec
SRM-FFT	149 sec	0.5 sec
FM-FFT	163,095 sec	0.5 sec

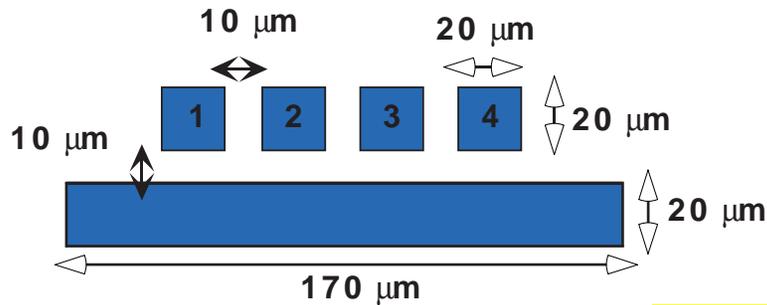
# Example II: method of characteristics (MC)

- Branin (1967) for single lossless line, Ho (1973) for multi-lossless lines, Gruodis (1979) for multi-resistive lines(R-L-C)
- objective: estimate waveform by applying Ho's method to lossless-like equation with similar efficiency as R-L-C circuit analysis
- represent waveform as combination of incident and reflected wave
- needs additional segmentation to include effect of distributed resistance
- additional voltage source needed for skin effect model



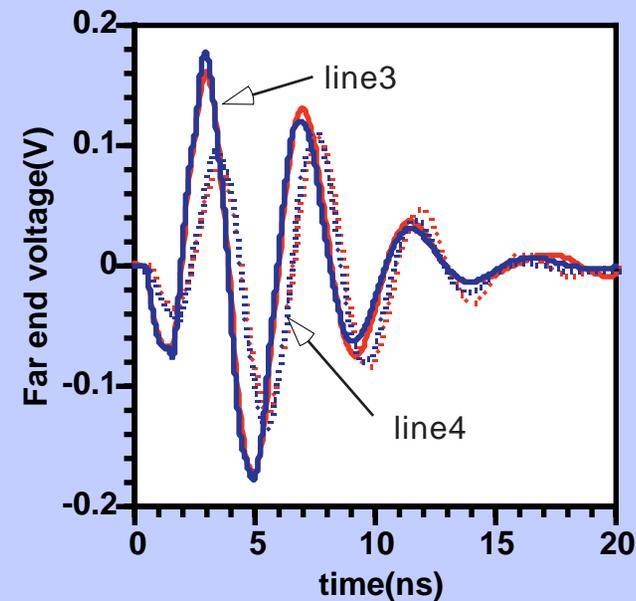
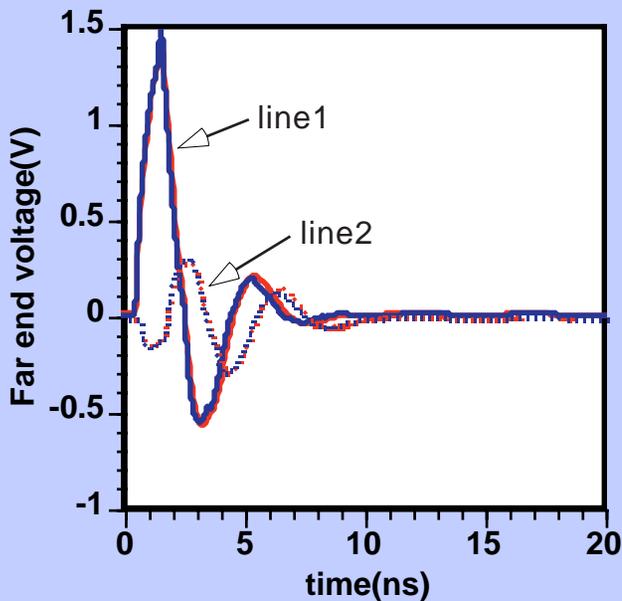
$$\begin{aligned}
 [V_R]_k^n &= [V]_{k+1}^{n-1} - [Z_o] \cdot [I]_{k+1}^{n-1} \\
 [V_I]_k^n &= [V]_{k-1}^{n-1} + [Z_o] \cdot [I]_{k-1}^{n-1} \\
 [Z_o] &= v_p \cdot [L'] \\
 [E_R]_k^n &= \Delta t \cdot v_p [V_{ds}]_{k+1}^{n-1} \\
 [E_I]_k^n &= \Delta t \cdot v_p [V_{ds}]_{k-1}^{n-1}
 \end{aligned}$$

# Four lines via method of characteristics



- line 1 excited,  $R_S=5 \text{ Ohm}$ ,  $C_L=10 \text{ pF}$
- run time comparison on Pentium II
  - SRM + MC: 7.6 sec
  - SRM + FFT: 283 sec

blue: FFT, red: MC



# Efficient interconnect modeling from dc to the skin effect

---

- **finite conductivity produces frequency dependent inductance and resistance**
  - important for accurate loss or cross-talk modeling
- **effective internal impedance and the surface ribbon method for dispersive R & L calculation**
  - excellent approximation from dc to high frequency
  - numerically very efficient
  - applicable to n-conductor and 3-D systems
- **small, frequency-independent R-L ladders can provide excellent equivalent circuit for frequency dependencies**
- **very efficient time domain conductor boundary condition demonstrated**