Copyright by

Olin Lee Hartin

Quantum Transport Simulations of Novel Compound Semiconductor Devices

by

Olin Lee Hartin, B.S., M.S.

Dissertation

Presented to the Faculty of the Graduate School of The University of Texas at Austin

in Partial Fulfillment

of the Requirements

for the Degree of

Doctor of Philosophy

The University of Texas at Austin May, 1998

Quantum Transport Simulations of Novel Compound Semiconductor Devices

Approved by Dissertation Committee:

Supervisor

Acknowledgements

I would like to acknowledge the help and support of all Team Neikirk members, especially Kiran Gullapalli. I must also thank the MBE group for growing wafers as well as their help in fabrication. I thank Dr. Maziar for her advice and tutelage on simulation matters, and Dr. Neikirk for his patience. Lastly, I thank my wife without whose help this dissertation would have been longer.

Quantum Transport Simulations of Novel Compound Semiconductor Devices

Publication No._____

Olin Lee Hartin, Ph.D. The University of Texas at Austin, 1998

Supervisor: Dean P. Neikirk

As devices in the semiconductor industry tend to shrink below 0.1 μ m quantum devices that work because of their small size, rather than in spite of it, become more attractive. It may be useful to simulate the operation of these devices whose behavior depends upon quantum tunneling and interference effects using comprehensive simulation tools.

In this work a two dimensional Schrodinger Poisson self-consistent simulator is described and demonstrated. Multi-valley coupling of effective mass equations is demonstrated and evaluated. A one dimensional Schrodinger Poisson self-consistent algorithm based on the tight binding formalism is also described and applied to heterostructure devices. Data from simulations based on these methods are compared with experimental data..

The methods developed allow the study of devices exhibiting quantum coherence effects combined with space charge effects in the presence of complex band structures and high electric fields. Such characteristics are present in a variety of heterobarrier problems and in structures with ultra-thin oxides. Our selfconsistent tight binding algorithm has been tested on several device structures.

Table of Contents

TABLE OF FIGURES IX
TABLE OF TABLESXV
CHAPTER 1 INTRODUCTION 1
CHAPTER 2 QUANTUM SWITCHING
2 .1 MOTIVATION
2 .2 DEFINITION OF QUANTUM STORAGE
2 .3 Cellular Automaton
2 .4 RTD Based Logic
2 .5 Memory Switching12
2 .6 QUANTUM STORAGE DEVICE
2.7 SUMMARY18
CHAPTER 3 EFFECTIVE MASS APPROXIMATION20
3 .1 BACKGROUND
3 .2 GREEN'S FUNCTION
3 .3 TIME INDEPENDENT EFFECTIVE MASS EQUATION
3 .4 2D DISCRETIZATION
3 .5 Homogeneous Solution
3 .6 INHOMOGENEOUS SOLUTION
3 .7 Concentration Calculation
3 .8 CURRENT CALCULATION
3 .9 Tests of the Algorithm
3.9.1 One Dimensional Simulation
3 .9.2 Two Dimensional Simulation
3.10 Summary
CHAPTER 4 TIGHT BINDING APPROXIMATION51
4 .1 Background
4 .2 Tight Binding

4.3 BAND STRUCTURES	62
4 .4 DISCRETIZATION	63
4 .5 Transfer Matrix method	66
4 .6 QUANTUM TRANSMITTING BOUNDARY METHOD (QTBM)	69
4.7 Self-Consistent Simulation	
4.8 Concentration	70
4.9 Current	72
4 .10 Results	73
4 .11 Summary	
CHAPTER 5 MULTI VALLEY EFFECTIVE MASS APPROXIMATION.	89
5 .1 Motivation	
5.2 Multi-valley	
5.3 SUMMARY	96
CHAPTER 6 QUANTUM STORAGE DEVICES	98
6 .1 BACKGROUND	
6 .2 Memory Switching Phenomena in Quantum Well Diode	100
6.3 QSD modeling and device physics	103
6 .4 Three terminal multi-state Quantum Storage Devic	
6 .5 2D SIMULATIONS OF QSD'S	
6.6 SUMMARY	
Appendix A	
Appendix B	
Appendix C	
Bibliography	
VITA	

Table of Figures

Figure 2.1: A QCA wire is shown where the charge state at one end of the array of dots effects the charge distribution at the other. Here dark dots contain charge and clear dots do not. Coulombic forces cause charge to align as shown. In referring to occupied dots the numbering scheme shown is used	
Figure 2.2	
Figure 2.3: This is a typical current density versus bias curve for a Double Barrier Resonant Tunneling Diode (DBRTD). Here a load line is shown as well. This is not quantitatively the load line used in this measurement 1310	
Figure 2.4: This is a memory cell based upon a RTD using load line switching ¹⁵ .	
Figure 2.5: These curves show several read write cycles of a QSD. The curves are grouped into states "1" and "2". Application of about 1.2 volts switches the device from curve "1" to curve "2". Application of about -1.2 volts switches the device from curve "2" to curve "1"	
Figure 3.1: This is the two dimensional discretization scheme. dz and dy are node spacings in z and y, respectively. The model space is indexed in i along z and j along j. $z_{i,j}$ is a solution at the node location (i,j)	
Figure 3.2: This is the symmetric Lanczos algorithm ³⁴ 26	
Figure 3.3: This is a flow chart of the process used to determine eigenvalue and eigenvectors	
Figure 3.4: This is the sparse matrix element structure	
Figure 3.5: This is the density of states (DOS) and transmission coefficint spectrum (t) at several locations in the DBRTD (Double Barrier Resonant Tunneling Diode) device shown above the graph. Curve 1 corresponds the beginning of the device at the contact, curve 2 corresponds to the end of the N+ region, curve 3 corresponds to the N- region adjacent to the barrier and curve 4 corresponds to the heterostructure quantum well. Note that the transmission coefficient in curve 5 peaks at about 0.2 eV. This coincides with the peak in the DOS spectrum of curve 4 which is the heterostructure quantum well. All other curves show a minimum at this energy indicating the electron lifetime is small except in the well. The other maxima and minima particularly in curve 1 are due to interference between incident wave and the	

Figure 3.6: On the left is a self-consistent solver flow chart and on the right is an illustration of the convergence of the space charge and maximum potential update

- Figure 3.7: The electron concentration profile of a wide DBRTD. Here the concentration on either end is in the contact region and in between concentration is in the heterostructure quantum well. This is a wide model with 565Å between nodes. The solution is similar to independent solutions at 565Å spacing across the device46
- Figure 3.9 This is the concentration profile in a very narrow DBTRD. A barrier is used on the sides to simulate Fermi level pinning. The high concentration on either end is in the contact region. The N++ regions show lateral interference effects.......48

- Figure 4.11: This figure shows the potential and concentration profile for this DBRTD. The solid curve is the tight binding approximation and the dashed curve is the effective mass approximation. Note that the concentration is very similar except in the barrier region where the tight binding concentration is larger, as expected. As a consequence the potential profile from the tight binding simulation is about 13% larger...82
- Figure 4.12: These are the tight binding simulation potential and concentration profiles at zero and 0.30 volts bias. Note the upwind potential barrier at position 40...82
- Figure 4.14: This is a plot of current density versus bias voltage for this DBRTD using several assumptions. Curve 1 is a self-consistent simulation based on the effective mass approximation. Curve 2 is a non self-consistent tight binding simulation assuming a straight line potential approximation. Curve 3 is a non self-consistent tight

Figure 4.16: The potential profile is shown to the left and the concentration profile is shown to the right. Two curves are shown. The dashed one is a Thomas Fermi simulation and the solid one is a tight binding simulation. Note the interference minimum that is located in the vacinity of the pulse doped region.....

- Figure 5.2: This is the potential profile of the coupled and tight binding simulations. The arrow shows the two best matches where the dark curve is the tight binding simulation and the light curve is the coupled effective mass simulation with $S_{\Gamma,X} = 0.42.92$
- Figure 5.3: This is a comparison between the tight binding and coupled effective mass simulations for an AlAs barrier of 17Å......93
- Figure 5.4: Using the coupling parameter $S_{\Gamma,X} = 0.35$ the concentration in the barrier region is a fairly good match. In the heterostructure itself the concentration is flat.

- Figure 6.1: A QSD structure is shown on the left and currents from Schrödinger Poisson self-consistent simulations of this structure are shown on the right.

- Figure 6.5: This is a suggested device structure with 100Å N++ region...... 107

- Figure 6.9: The graph on the left shows the density of states (DOS) spectrum for the N++ layer and the heterostructure quantum well (HQW). The transmission coefficient (t) spectra is also shown. The transmission resonance peak is at about 0.18 eV for the solution on the left and 0.16 eV for the solution on the right. These correspond to the left and right solutions above, respectively. At these transmission resonances there is a DOS node elsewhere in the device. The resonance at 0.038 eV in the solution on the

left peaks in the N++ regions and is diminished but still a resonance elsewhere. Resonances also occur at about 0.078 eV in the left solution and at about 0.08 eV on the right solution and are diminished in the heterostructure quantum well. . 112

- Figure 6.10: Γ , X, and L log concentration profiles are shown. Using this coupling
- Figure 6.11: This is the tight binding DOS and transmission spectra for a tight binding simulation. These curves are very similar to those shown with the effective mass approximation in Figure 6.9. Transmission coefficients are generally higher at low energy than with the effective mass approximation. The solution on the left has similar resonances except that the HQW peak is at about 0.2 eV which is about 0.02ev above the corresponding solution in Figure 6.9. This is due to the higher potential in that portion of the device for the tight binding solution. The resonances on the right are shifted down about 0.015 eV from those on the right side in Figure 6.9.114

Figure 6.12: This is a memory circuit implemented with a three terminal QSD (TQSD). 63

- Figure 6.13: This plot shows the advantage of on/off resistance ratio in the performance of
- Figure 6.15: This is a triple barrier device structure. It is similar to the read terminal
- Figure 6.16: These are current ratios between solutions for a range of device structure parameters. This suggests AlAs barriers of 17Å and QW width of 50Å or 100 Å.
- Figure 6.17: On the left are potential profiles from self-consistent Schrödinger Poisson solutions and on the right are concentration profiles. Curve 1 is a solution with quasibound state solutions and curve 2 is a solution without quasi-bound state solutions. The top two graphs are based on the effective mass approximation and the bottom two
- Figure 6.18: The current density versus bias voltage plot shows two solutions. The inset region on the box in the curve on the left is shown in the curve on the right. Curve 1 is the quasi-bound state solution, curve 2 is the solution not supporting quasi-bound
- Figure 6.19: Current density versus bias curves measured in the laboratory show two solutions. Both solutions show inflections suggesting resonance peaks and valleys.

Table of Tables

 Table 4.2: The valence band offset at the GaAs/AlAs heterostructure interface at room temperature using the simple ratio.

 61

Chapter 1 Introduction

There are a number of novel devices that depend upon quantum tunneling and interference effects. Since some of these device ideas are difficult to test in the laboratory, the need to do optimization and inverse modeling in design of these devices suggests development of more comprehensive simulation tools.

Effective mass approximation-based Schrödinger Poisson simulation tools make it possible to rapidly simulate large device models. Convergence is an issue in part because the density of states function is highly nonlinear in these problems. The tight binding Hamiltonian can be used to do simulations of a range of materials including band mixing between materials. Valley mixing affects carrier concentration and transmission in devices with complex structures. Less rigorous methods based on effective mass approximations may be used to approximate these effects. However, there are differences between simulations based on coupling of effective mass equations and on the tight binding approximation.

One class of novel devices is simulated with these methods. The quantum storage device (QSD) is one of the new class of novel devices based on simulations and laboratory measurements. The addition of simulation methods introduced here add to the understanding of this device. Self-consistent solutions to the Schrödinger and Poisson equations have been widely used to identify both qualitative, and with varying degrees of success, quantitative behavior of Double Barrier Resonant Tunneling Diodes (DBRTDs).^{1,2}. Self-consistent solutions are essential because quantum well diodes often incorporate lightly doped layers, and

the resulting space-charge effects can significantly influence device $characteristics^{3,4}$.